



DE LA RECHERCHE À L'INDUSTRIE

# Adaptive Multilevel Splitting for Particle Transport

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Commissariat à l'énergie atomique et aux énergies alternatives - [www.cea.fr](http://www.cea.fr)

## ❑ Collaborative work between :

- E. Dumonteil (CEA/DRF/IRFU/DPhN)
- K. Frohlicher (IRSN/PSN-RES/SNC) → PhD 2022
- T. Lelievre (Ecole des Ponts/CERMICS)
- H. Louvin (CEA/DRF/IRFU/DEDIP) → PhD 2017
- D. Mancusi (CEA/DES/ISAS/DM2S)
- M. Nowak (Previsions.io) → PhD 2018
- M. Rousset (INRIA Rennes)
- L. Thulliez (CEA/DRF/IRFU/DPhN)

Past & ongoing publications

- Louvin et al, EPJ Nuclear Sci. Technol. Vol 3 (2017)
- Nowak et al, Nuc. Sci. Eng., Vol 193 (2019)
- Mancusi et al, Trans. Am. Nuc. Soc., Vol. 120 (2019)
- Frohlicher et al, ongoing work (2021)
- Thulliez et al, ongoing work (2021)

## □ Context

- Motivations
- AMS for MC particle transport codes: basic ideas
- Implementation & validation

## □ Adaptation of AMS for particle transport

- On-the-fly scoring
- Branching tracks
- Multi-particles/particles cascades
- Towards self-learning of the cost function

## □ Use of AMS in reactor physics

- Chain reaction & population control
- Spatial correlations
- AMS & branchless collisions
- Results

# □ Context

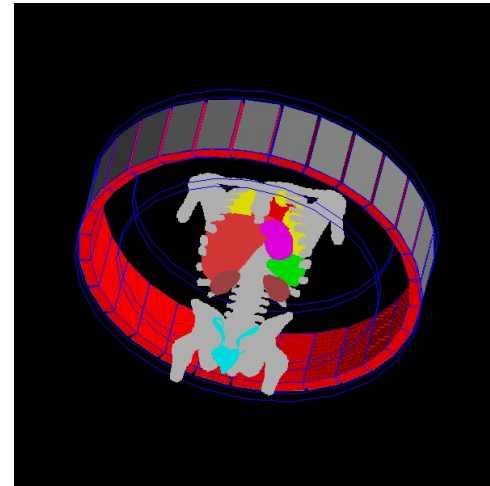
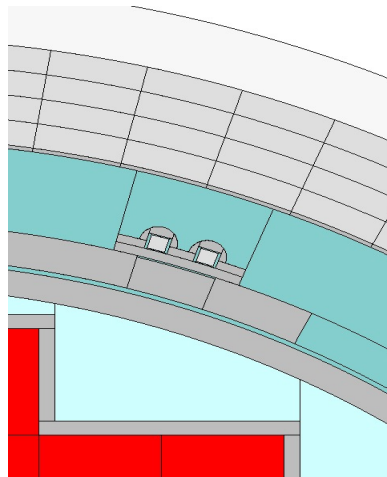
- Motivations
- AMS for MC particle transport codes: basic ideas
- Implementation & validation

- ❑ Particle physics in its large sense: radiation protection, medical, fundamental, etc.
- ❑ **Boltzmann equation** in non-reproductive media and fixed sources
- ❑ Neutrons, photons, électrons, muons, etc. at all energies
- ❑ Main challenge: **variance reduction** w.r.t.  $x$  and  $E$  parameters

Ex-core dosimetry  
&  
PWR vessel fluence

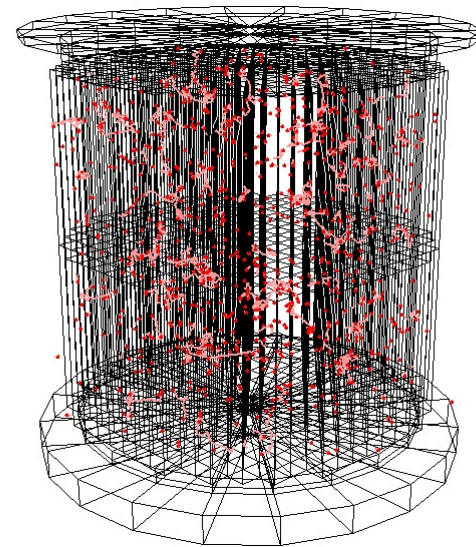
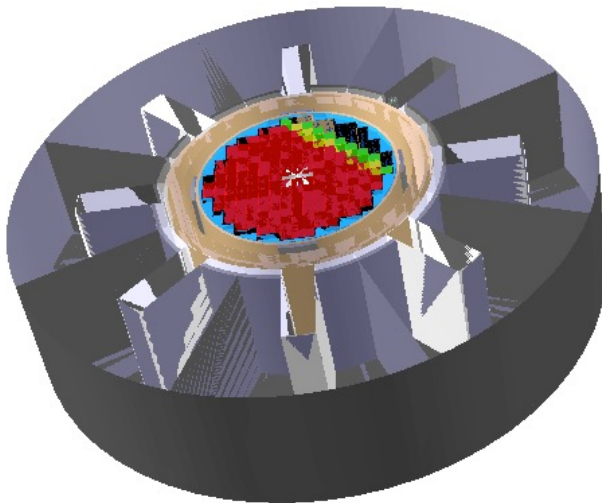


Radioactive waste  
management



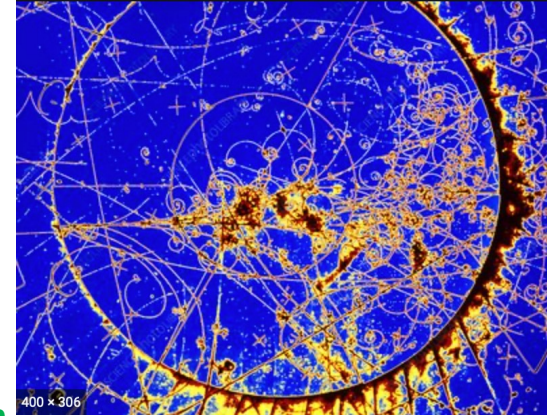
Radiotherapy  
([opengatecollaboration.org](http://opengatecollaboration.org))

- ❑ Static **linear Boltzmann equation** in fissile media (no fixed sources)
- ❑ Neutrons between 0 and 20 MeV
- ❑ Main challenge: finding **eigenvalues and eigenvectors** while tackling with correlations

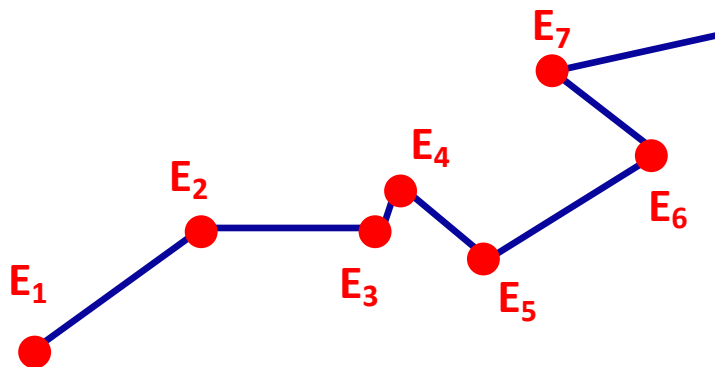


- ❑ **Originates** from applied mathematics applied to molecular dynamic
  - (Cerou et al, 2007)
  - (Cerou et al, 2011)
  - (Aristoff et al, 2015)
  
- ❑ **Adaptation to particle transport**
  - CEMRACS@CIRM 2013 (Lelievre & Dumonteil)
  - (Louvin, Dumonteil, Lelievre, Rousset 2017)
  - (Louvin, Mancusi & Dumonteil 2019)
  
- ❑ **Objective of the different developments presented is to fit in the framework of AMS for discrete Markov chains** detailed in (Brehier et al, 2016)

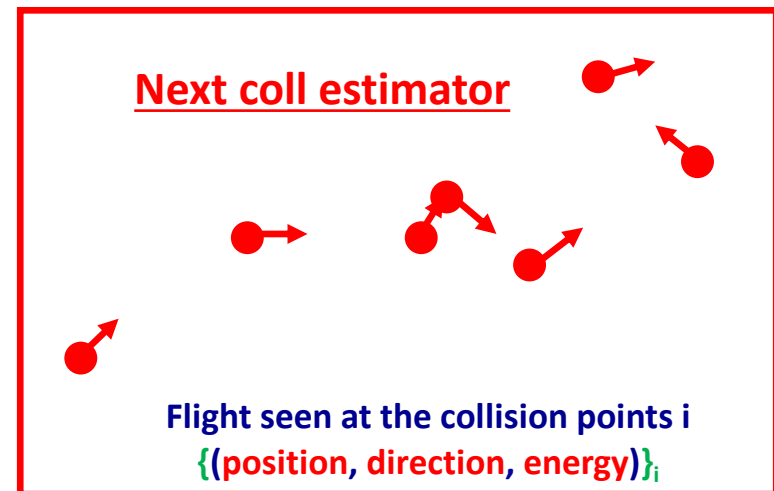
- ❑ Term "tracks" originates from first bubble chambers
- ❑ Parameter space (**position, direction, energy, time, particle type**)
- ❑ **Non-charged** particle tracks are
  - Markovian **everywhere** in (**position, direction, energy**)
  - Markovian at **collision points**  $i$   **$\{(\text{position}, \text{energy})\}_i$**
  - Markovian at **collision points**  $i$   **$\{(\text{position}, \text{direction}, \text{energy})\}_i$**
  - ...
- ❑ AMS for discrete Markov chains is the most suitable AMS « flavor »



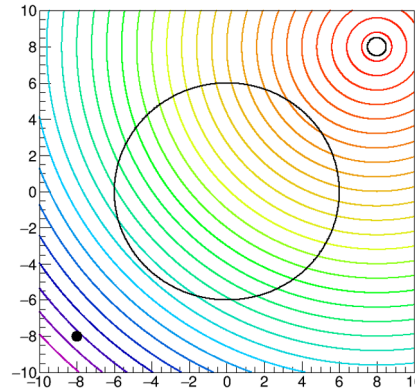
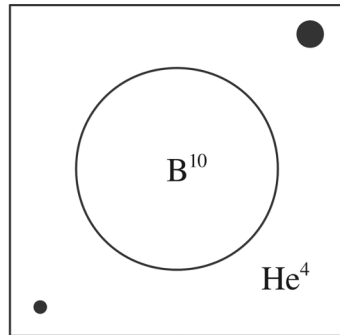
Tracks left by a neutrino interacting with a neutron



Typical exponential flight undergone by a neutron with anisotropic collision kernel  $C(E)$

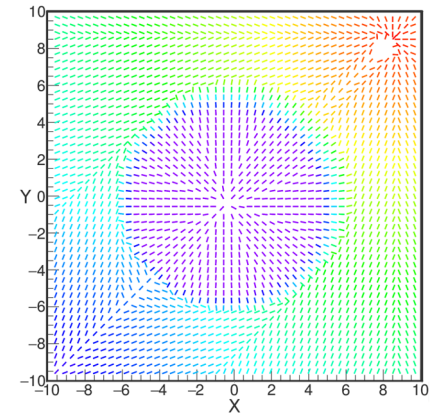




Example of cost functions:

Simple only-spatial  
 $\xi = 1/(1+d)$

Material+energy taken  
into account in  $\xi$

Algorithm with parameters

$\left\{ \begin{array}{l} n \text{ (\# of particles)} \\ k \text{ (\# of particles duplicated/iteration)} \\ \xi \text{ (cost function)} \end{array} \right.$

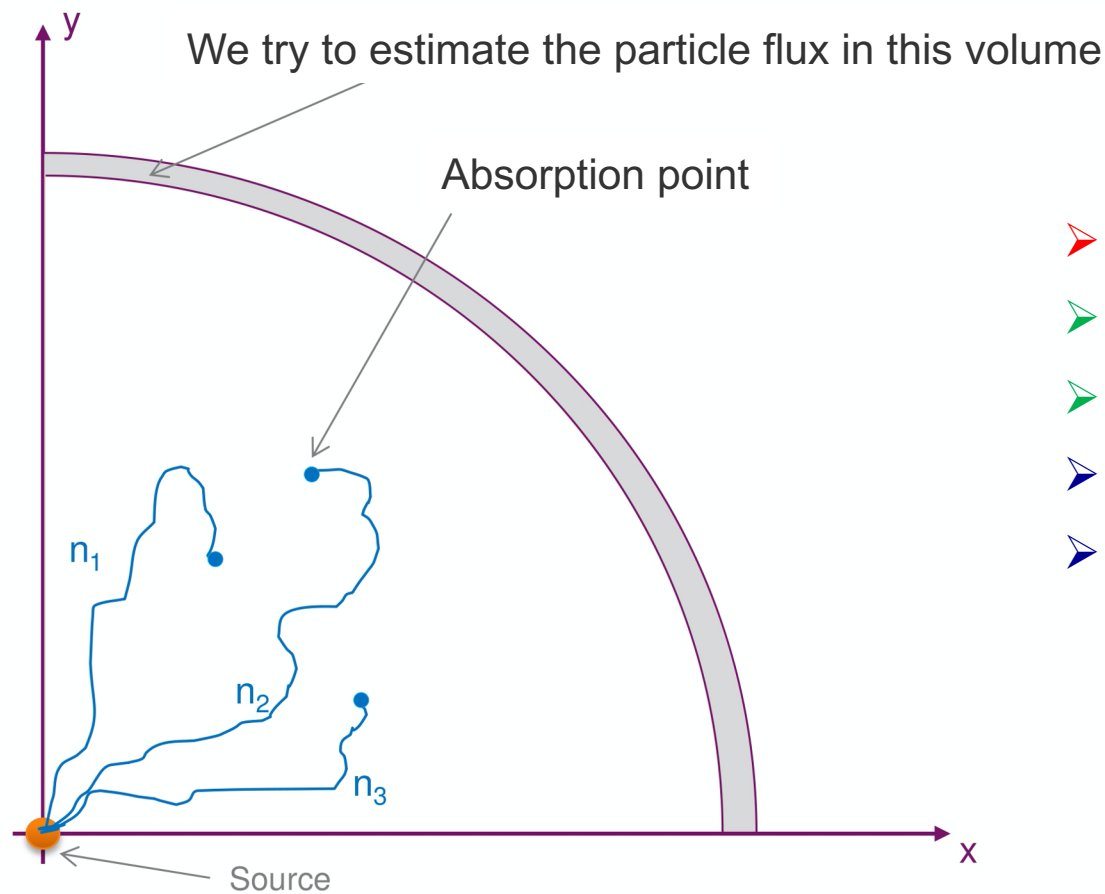
- $n$  particles simulated => tracks are stored
- score is assigned to each neutron track (= Max of  $\xi$  over whole trajectory)
- tracks are ranked according to their score
- the  $k$ -th “worst” track defines the new splitting level
- the  $k$  tracks having scores lesser than this level are deleted
- $k$  tracks are randomly selected and duplicated at the splitting level
- a new set of  $n$  particles is obtained, and we start the whole process again

□ **Stopping criterium:**

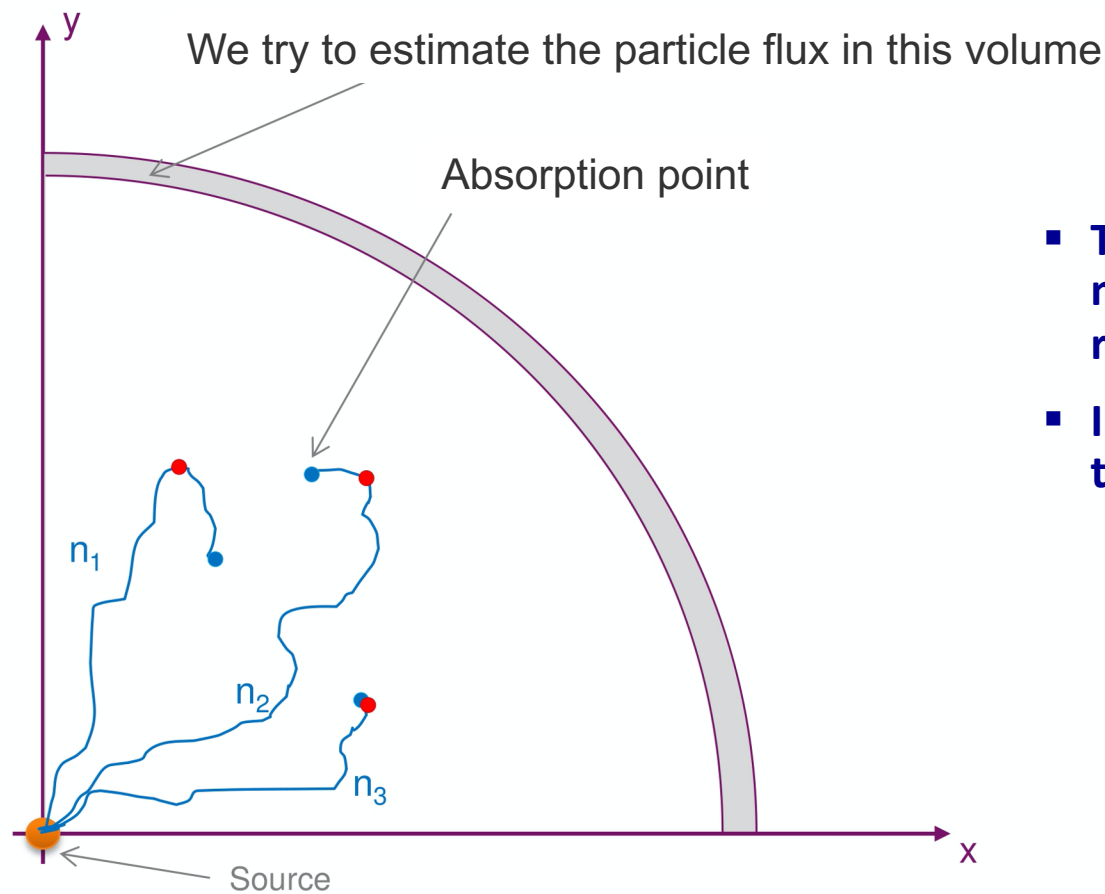
- **When n-k tracks have reached the “detector”, the algorithm stops**
- **The number of iteration corresponding to reach that criterium is N**
- **Each neutron is assigned a statistical weight  $\alpha$  being:**

$$\alpha = \left(1 - \frac{k}{n}\right)^N$$

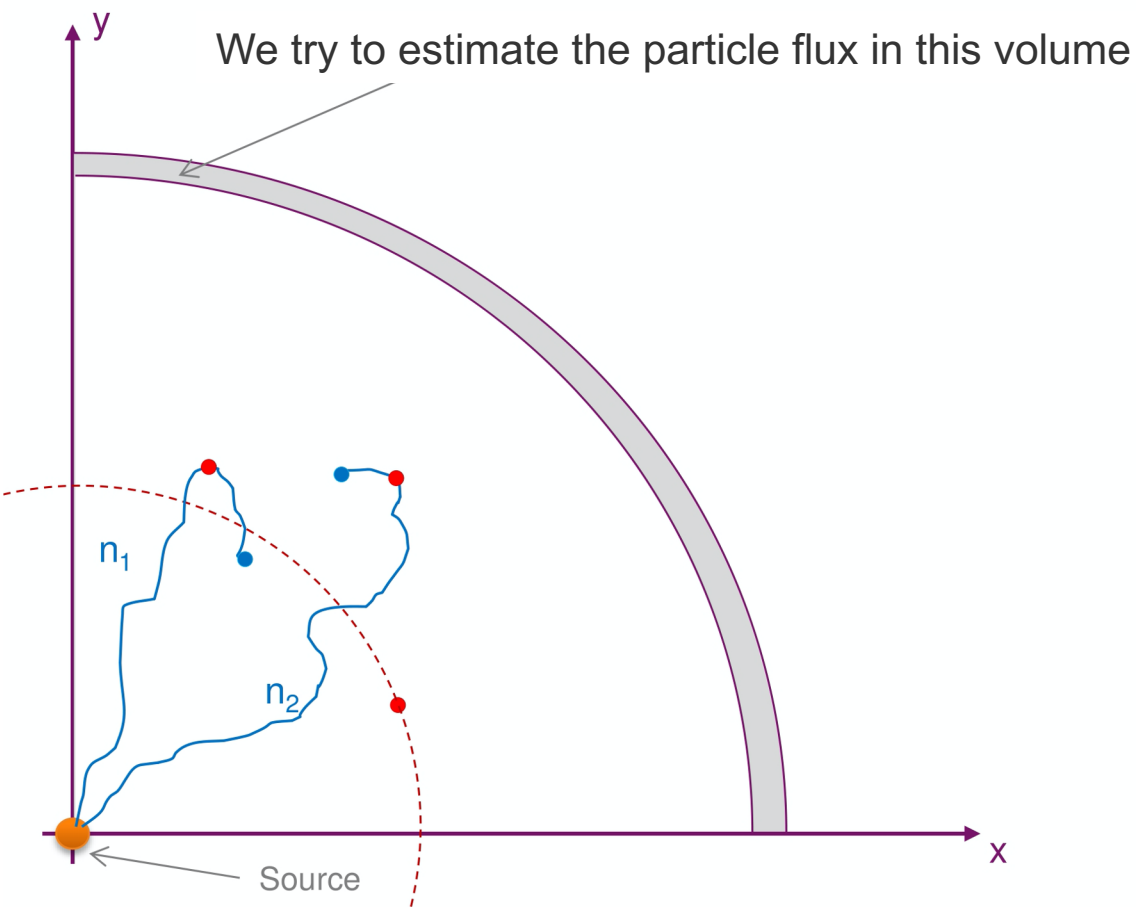
- **An unbiased estimator of the flux is calculated “as usual” using the tracks of the last iteration**



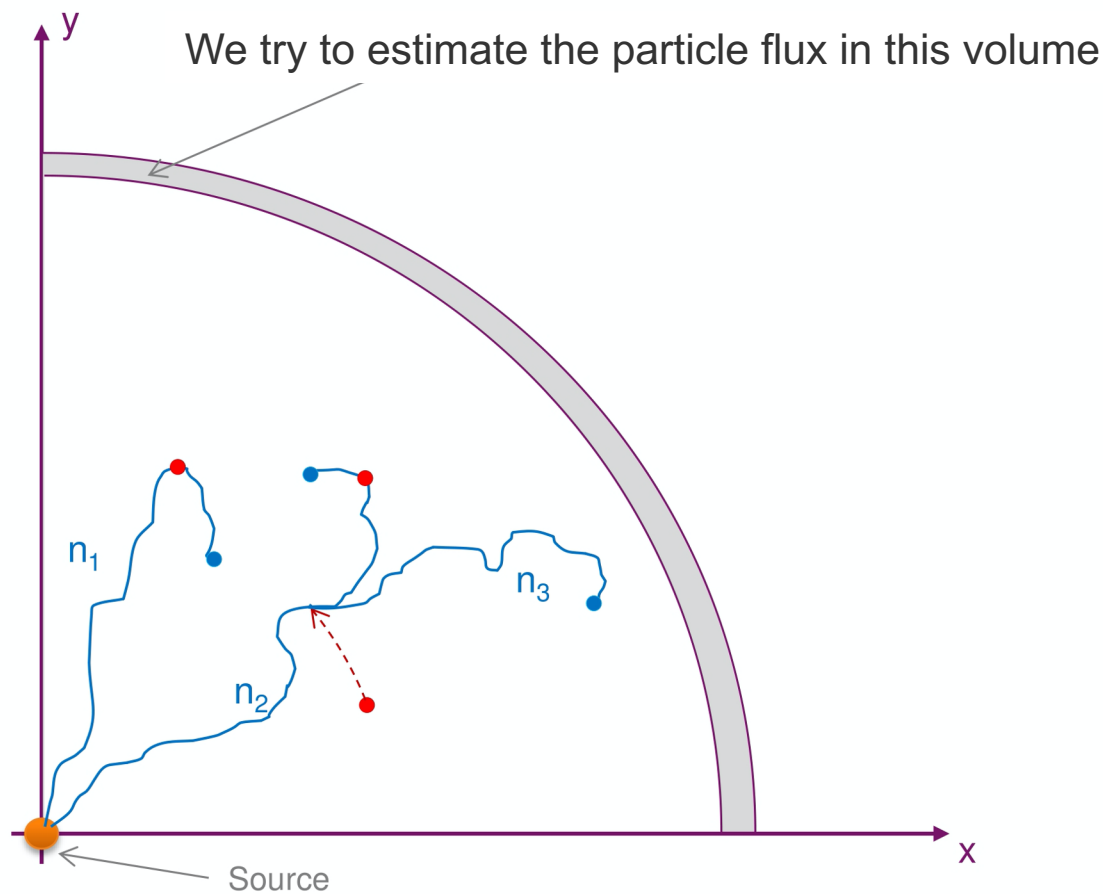
- $n=3$
- $k=1$
- $\xi$ : distance to the source
- Target: spherical shell (purple)
- 3 particles simulated from the source to their absorption (blue points)



- The importance function is the maximal distance to the source reached by the particle (red points)
- In this case the neutron tracks with the lowest score is  $n_3$

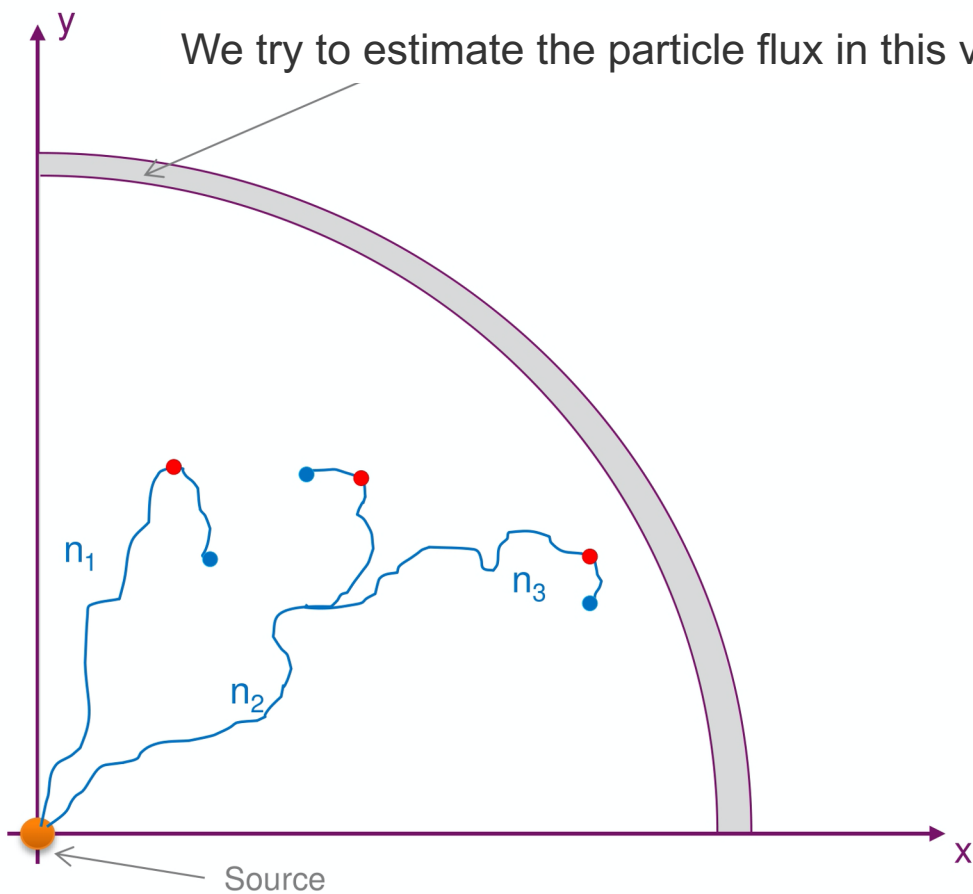


- The tracks associated to particle 3 is deleted
- The maximum score of this track is stored and defines the first splitting level



- Track number 2 is randomly sampled for the splitting
- A new particle is simulated from this splitting point until its absorption

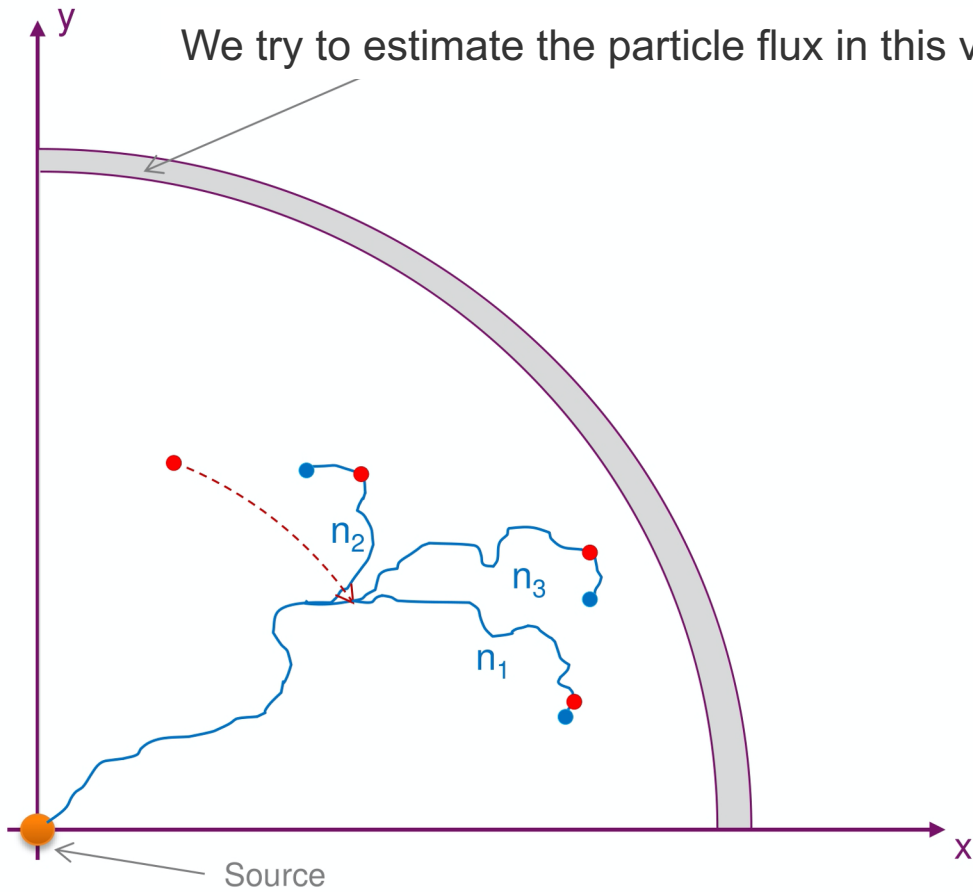
We try to estimate the particle flux in this volume



- The score of this new tracks  $n_3$  is calculated
- The first iteration is over
- The stopping criterium is not meet: the iteration process goes on

We try to estimate the particle flux in this volume

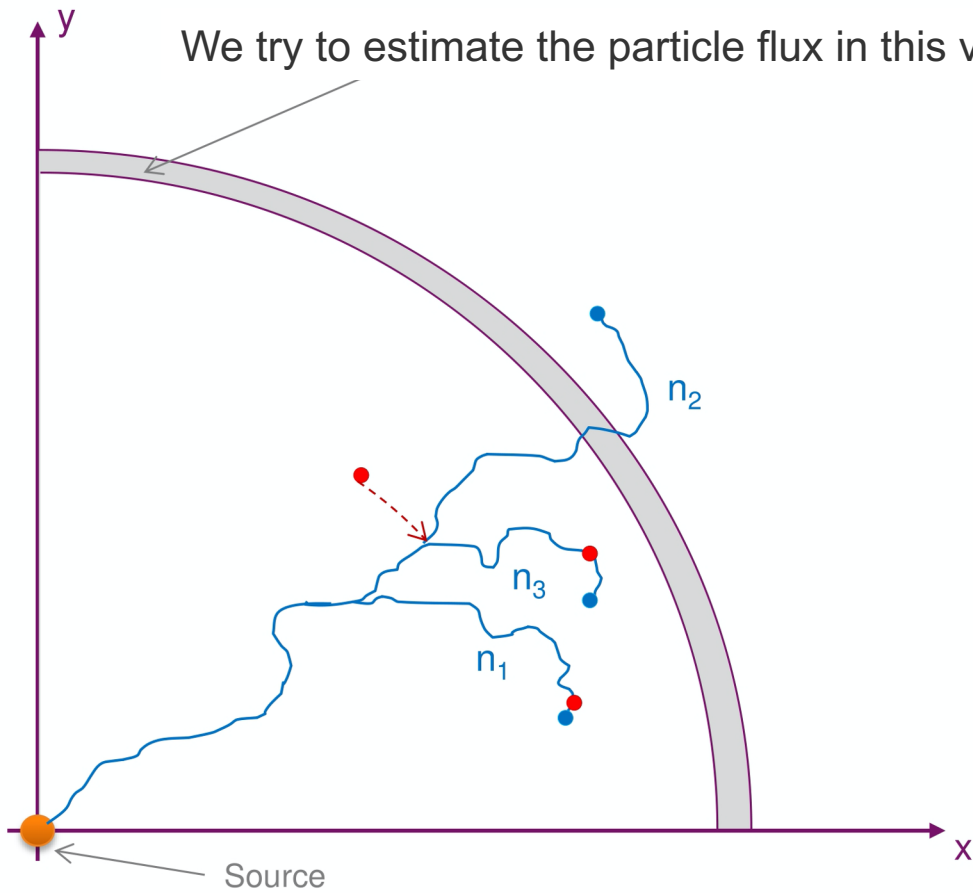
- **Iteration 2**





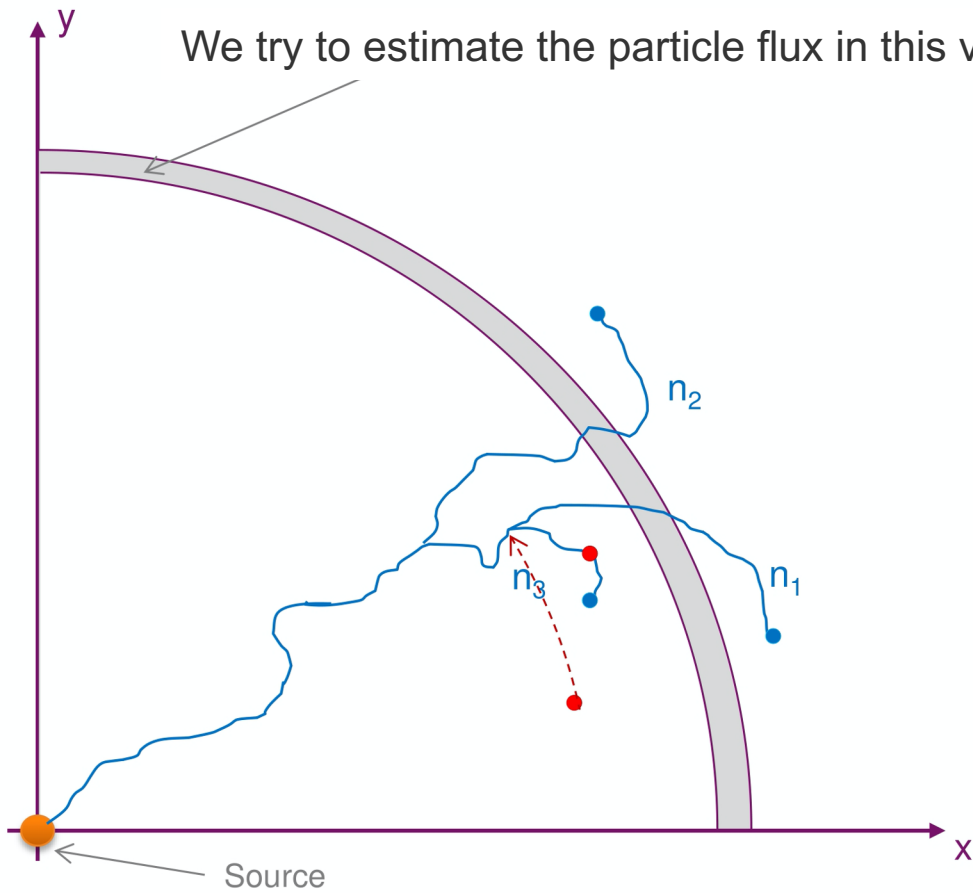
We try to estimate the particle flux in this volume

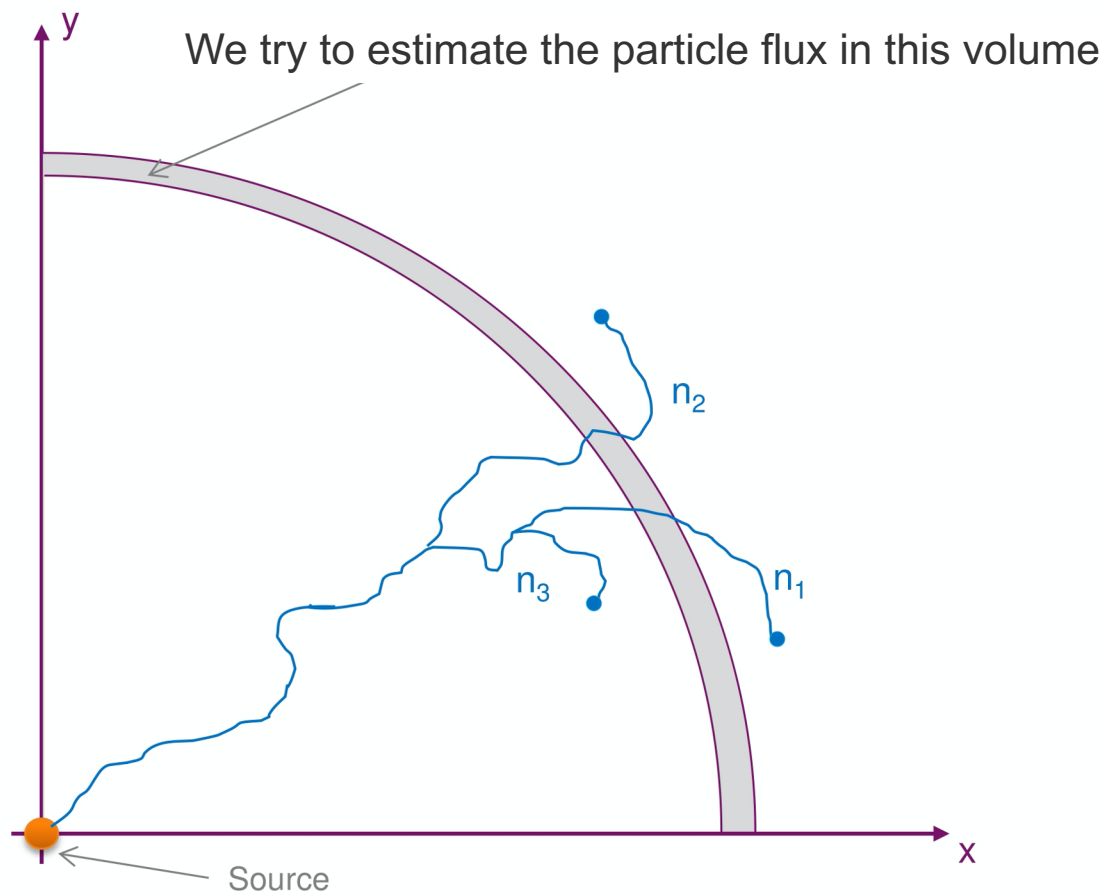
- Iteration 3



We try to estimate the particle flux in this volume

- Iteration 4

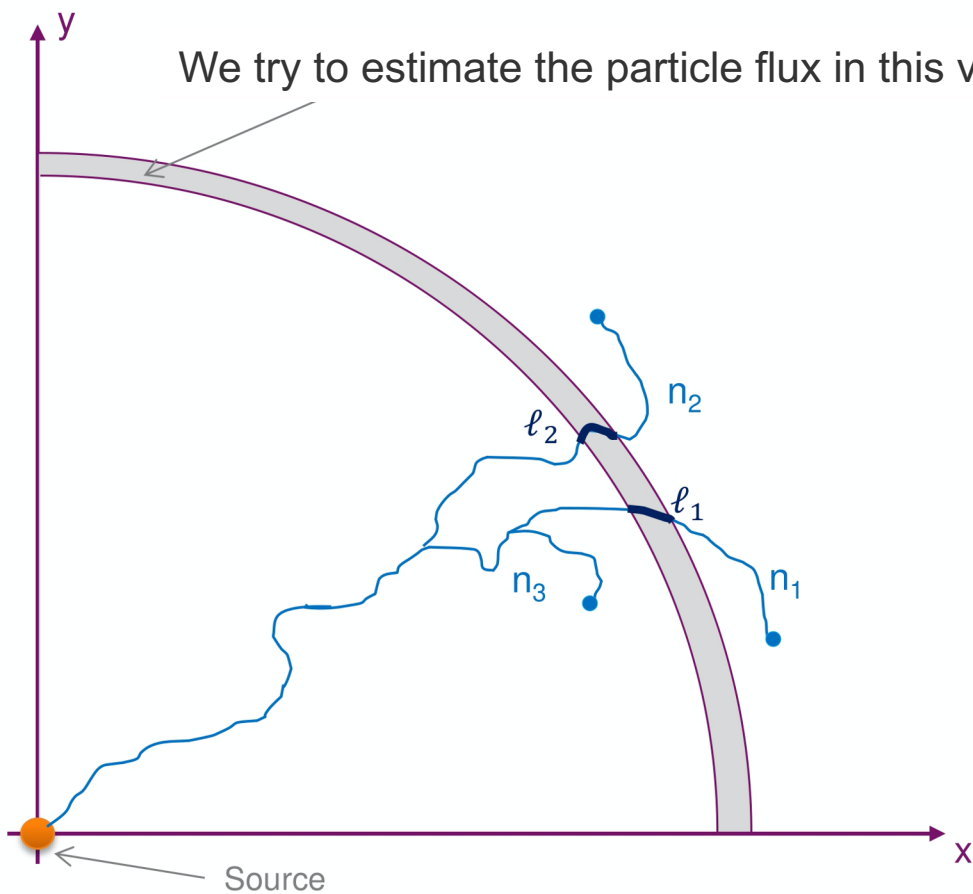




- Iteration 4
- **n-k** particles have reached the target, the algorithm stops
- The statistical weight of the particles is :

$$\alpha = \left(1 - \frac{1}{3}\right)^4$$

We try to estimate the particle flux in this volume



- The flux is calculated according to standard MC flux estimators. For example the travelled length in the spherical shell can be used to tally the flux:

$$\varphi = \frac{1}{3} \alpha (l_1 + l_2)$$



TRIPOLI 4®



- TRIPOLI-4 @ CEA
- Distributed by OECD/NEA
- Neutron, gamma, e+, e-
- E < 20 MeV
- Evaluated cross-sections
- (Brun et al, 2015)
- [www.cea.fr/energies/tripoli-4](http://www.cea.fr/energies/tripoli-4)
- Nuclear industry

- Geant4 @ CERN
- Open-source
- All particles
- All energies
- Both evaluated cross-sections & models
- (Agostinelli et al, 2003)
- [geant4.web.cern.ch](http://geant4.web.cern.ch)
- Fundamental / medical / spatial

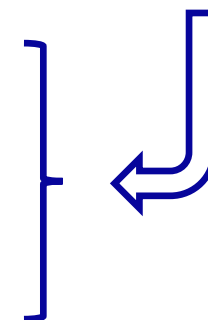
AMS has been implemented in a light (few classes) C++ framework

Will be released and distributed as open-source package in 2021

Small 'user guide' to link it to other transport codes

Provided with basic cost functions (distance to spatial detector, radial, ...)

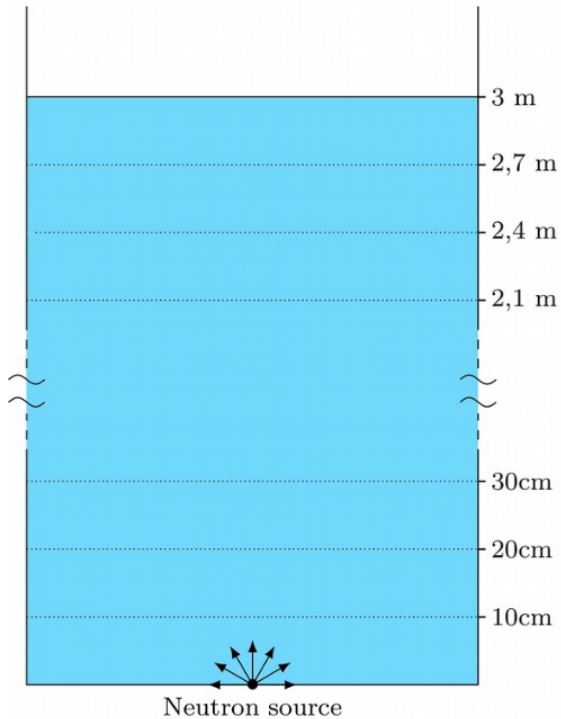
Verified through analytical benchmarks



# □ ADAPTATION OF AMS FOR PARTICLE TRANSPORT

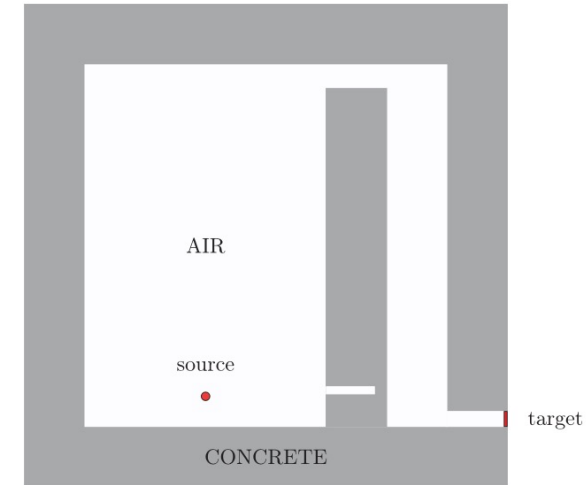
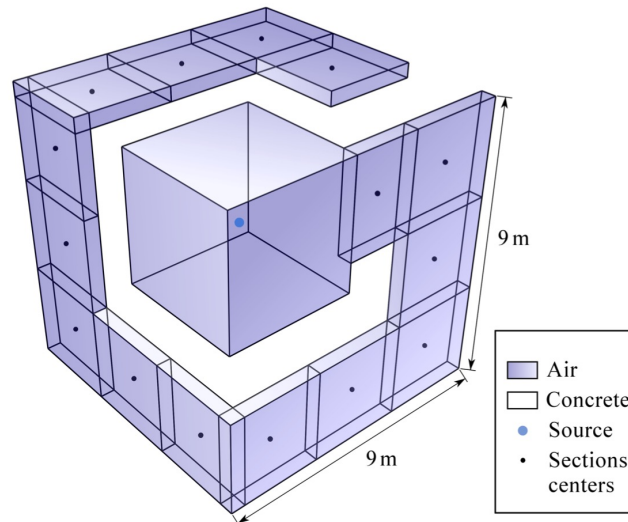
- On-the-fly scoring
- Branching tracks
- Multi-particles/particles cascades
- Towards self-learning of the cost function

- ☐ Many radiation protection problems require to score quantities "everywhere"



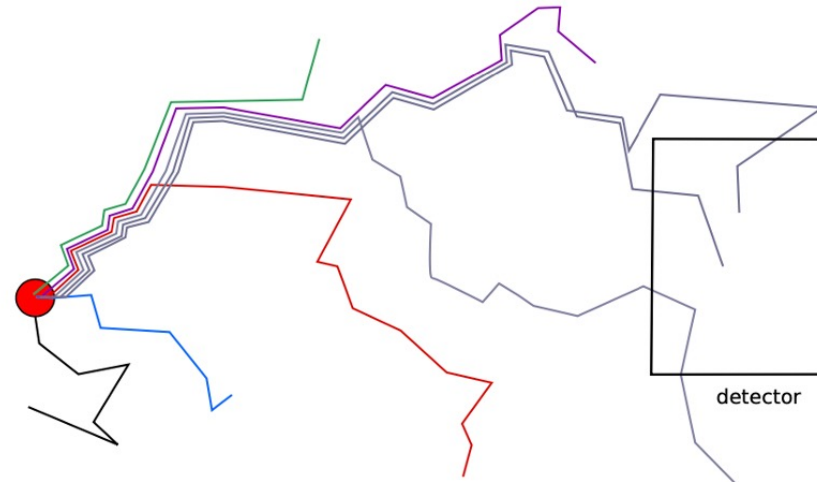
Flux calculation in the cooling pool of a reactor

Streaming problems as can be met in nuclear marine propulsion



Dose calculations while dimensioning shielding rooms

- Following (Brehier, 2016), the idea consists in building a **particle genealogy**
- **Old tracks** are kept in memory
- **New tracks** = copy of the track selected for duplication from 1st point up to splitting point



We consider the analog Monte Carlo **unbiased estimate** of a score  $\psi$ , where  $(X_i)_{i \in [1, n]}$  is a set of analog tracks

$$\hat{\psi}_{MC} = \frac{1}{n} \sum_{i=1}^n \psi(X_i),$$

For any iteration  $q$  of the scoring process we define the following **unbiased estimate**:

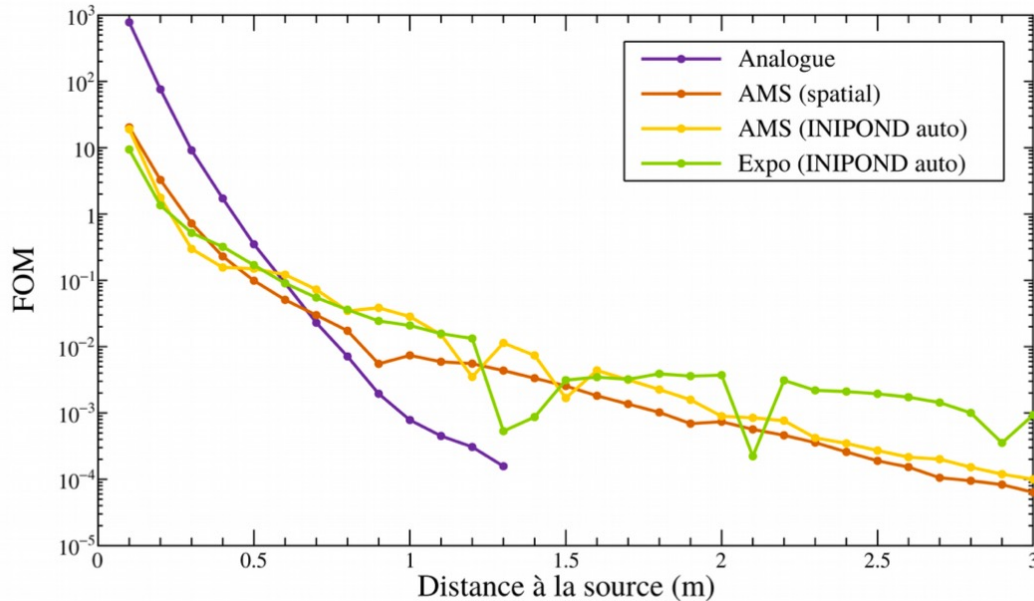
$$\hat{\psi}_q = w_q \hat{\psi}_q^{\text{on}} + \sum_{j=0}^q w_j \hat{\psi}_j^{\text{off}}, \quad \text{with } w_j = \begin{cases} \frac{1}{n} & \text{if } j = 0 \\ \frac{1}{n} \prod_{i=0}^{j-1} \left(1 - \frac{K_i}{n}\right) & \text{if } j > 0 \end{cases}$$

$$\hat{\psi}_q^{\text{on}} = \sum_{X \in T_q^{\text{on}} \rightarrow T/\xi > Z_q} \psi(X) \quad \text{and} \quad \hat{\psi}_q^{\text{off}} = \sum_{X \in T_q^{\text{off}} \rightarrow T/\xi < Z_q} \psi(X).$$

$K_q = \# \text{ tracks} / \xi < Z_q$   
 $Z_q = \text{splitting level}$



## ☐ Cooling water pool

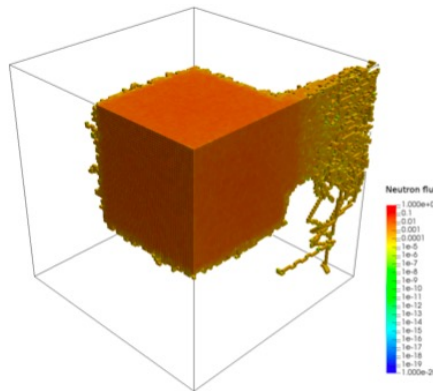


TRIPOLI-4 embedded  
deterministic solver for  
~adjoint flux calculation

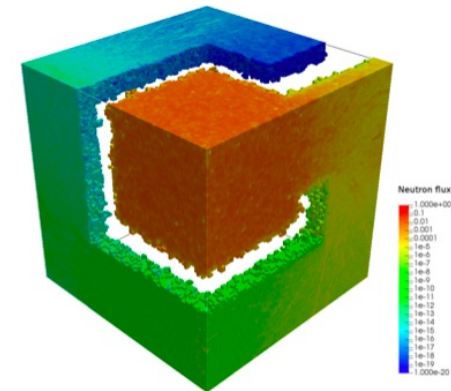
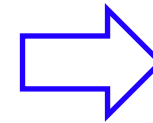
- $\xi$  either spatial or INIPOND
- AMS FOM close to ET ...
- ... even with naïve  $\xi$  !

## ☐ Labyrinth

- 2 MeV isotropic neutron source
- Air surrounded by concrete
- **ET fails** due to air concentration
- **AMS  $\xi = 0$**  in concrete, growing following path to exit



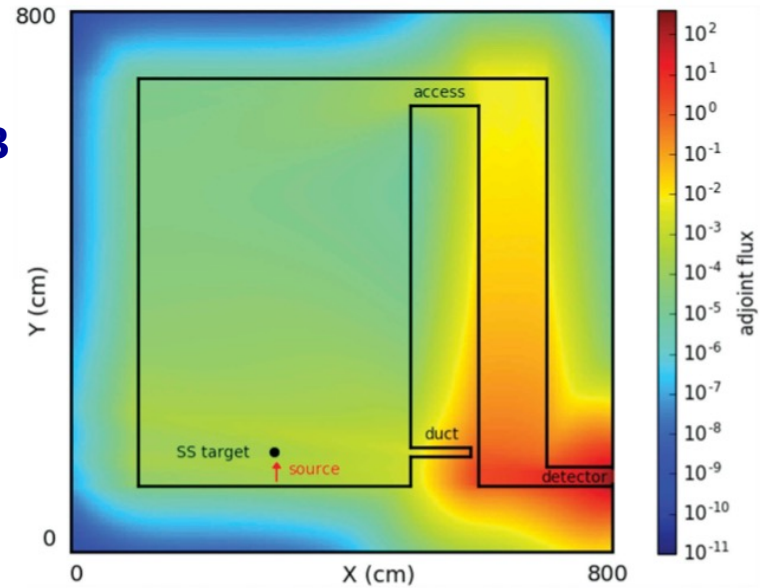
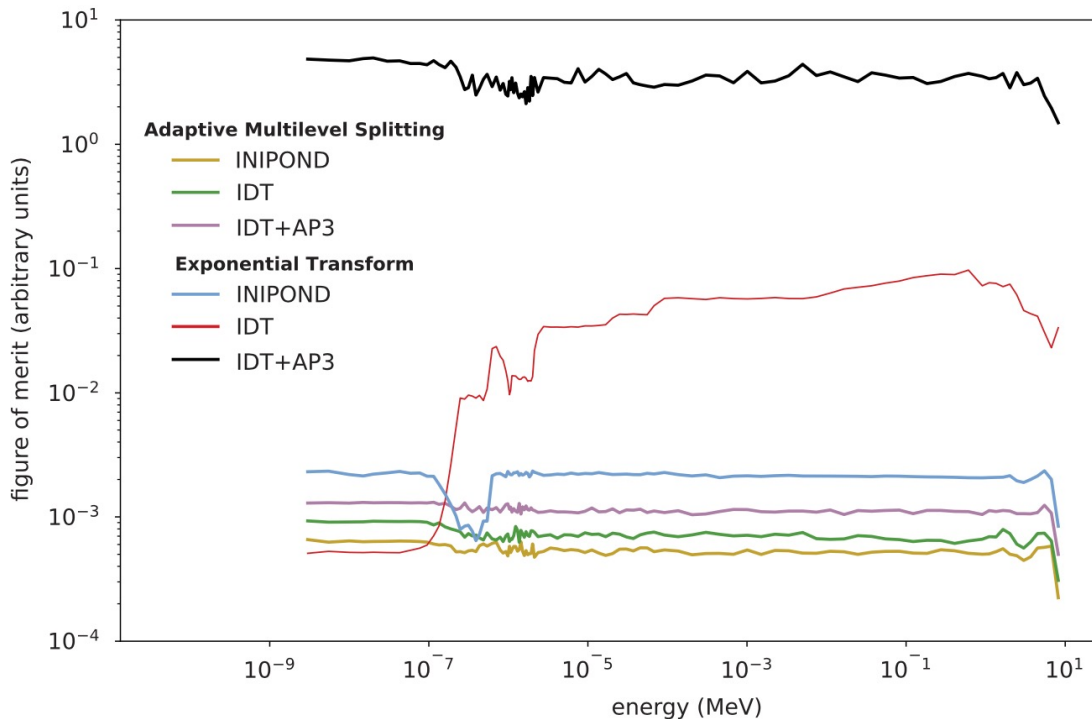
(a) Analog neutron flux



(b) AMS neutron flux

- ❑ **Bunker benchmark**
- ❑ **(Mancusi et al, 2018)**
- ❑ **First coupling between CADIS\* & AMS**
- ❑ **Comparison coarsest->finest: INIPOND/IDT/IDT+AP3**

$\xi$  = adjoint flux from IDT deterministic solver



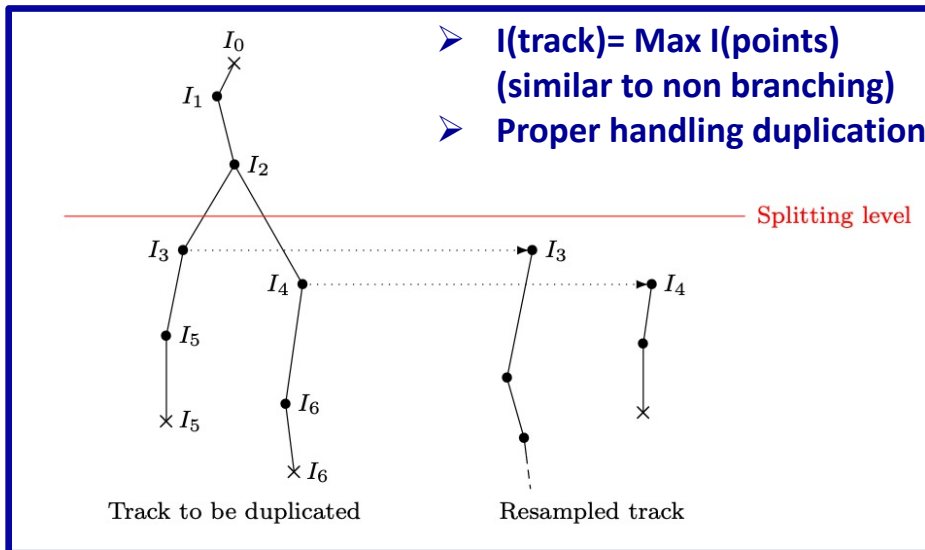
- **ET : larger FOM if  $\xi$  well-known**
- **AMS & ET exhibit similar FOM for 'intermediate'  $\xi$**

**CADIS\* = Consistent Adjoint Driven Importance Sampling (Haghighat, 2003)**

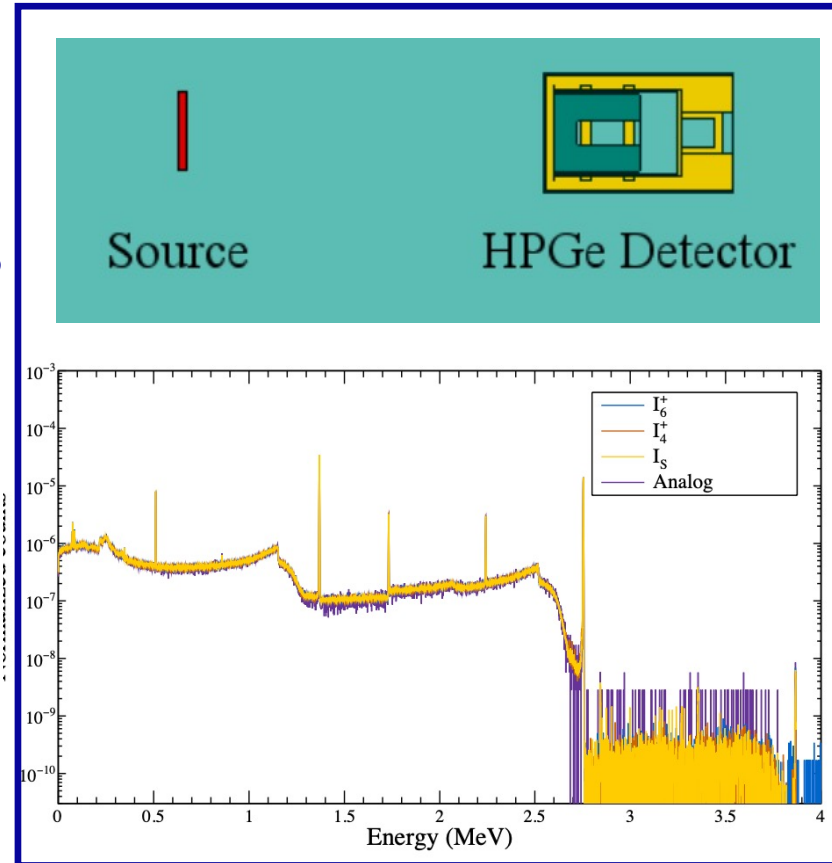
**IDT = 3D Cartesian deterministic solver for the multigroup time-independent transport equation (Zmijarevic, 2001)**

**AP3 = improved adjoint flux wrt anisotropy, upscattering, and energy during cross section condensation (Schneider, 2016)**

- ❑ Many examples of branching structures in particle transport:
  - (n,2n) ➤ electromagnetic cascade
  - fissions ➤ Intranuclear cascade
- ❑ Weights/particles could be introduced but correlations would be lost
- ❑ Handled by an appropriate algorithm within AMS so as to preserve correlations

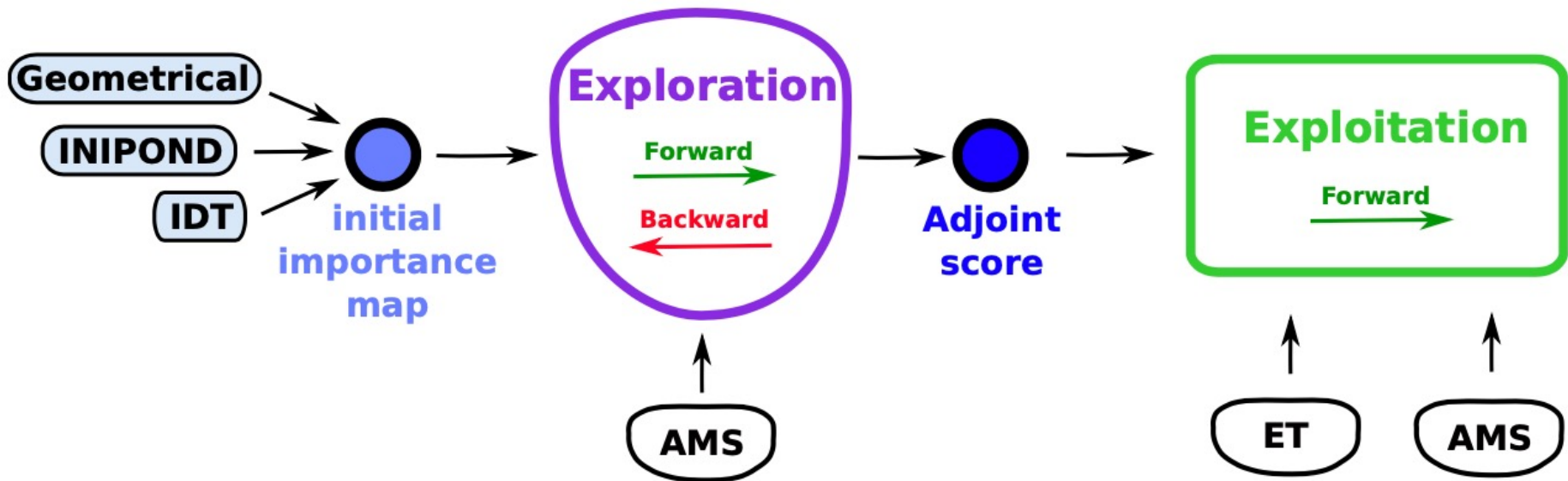


HPGe detector: g, e<sup>+</sup>, e<sup>-</sup>



➤  $FOM_{AMS}/FOM_{analog} \sim 10^2$

- Investigating the use of machine learning to improve the estimation of the adjoint flux
- (Nowak et al, 2018)
- Numerical 2-step scheme using AMS to adaptively improve the cost function / adjoint flux :



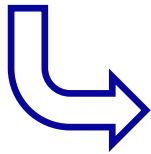
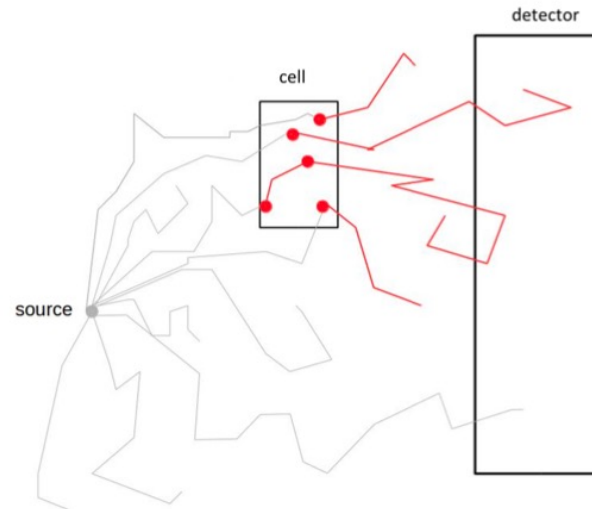
## Standard MC estimator of the adjoint flux

Expected contribution  $c$  of a point  
in the phase space

$$x = (\vec{r}, \vec{\Omega}, E)$$

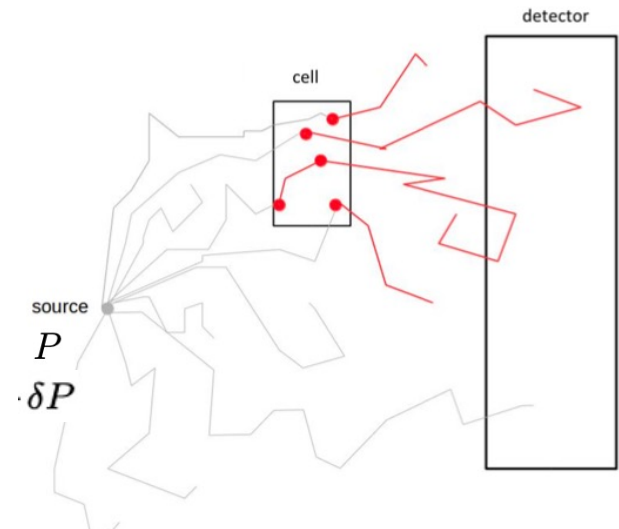
to a response (flux in a detector).

$$\psi^\dagger(P) = \mathbb{E}(c|P)$$



- $N$  histories  $\{T_0, \dots, T_N\}$
- $N_i$  points  $\{P_0^{(i)}, \dots, P_{N_i}^{(i)}\}$

$$\psi^\dagger(P) = \frac{\sum_{P_i \in \mathcal{T}} c_i \mathbb{1}_{\delta P}(\mathbf{r}_i, E_i, \Omega_i)}{\sum_{P_i \in \mathcal{T}} \mathbb{1}_{\delta P}(\mathbf{r}_i, E_i, \Omega_i)}$$



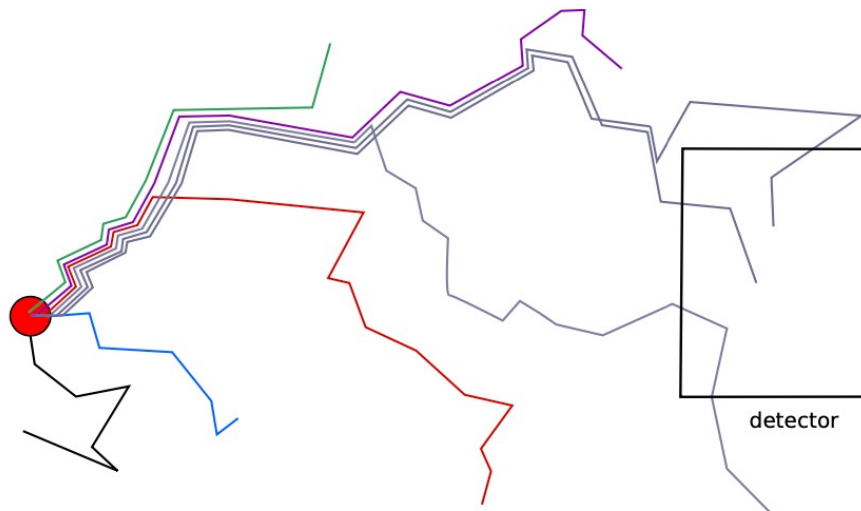
AMS estimator of the adjoint flux => the same with weights given by

$$\psi^\dagger(P) = \frac{\sum_{P_i \in \mathcal{T}} c_i \mathbf{1}_{\delta P}(\mathbf{r}_i, E_i, \Omega_i)}{\sum_{P_i \in \mathcal{T}} \mathbf{1}_{\delta P}(\mathbf{r}_i, E_i, \Omega_i)}$$

becomes

$$\frac{\sum_{\mathcal{T} \in T^{\text{AMS}}} w_{\text{AMS}}(\mathcal{T}) \sum_{P_i \in \mathcal{T}} c_i \mathbf{1}_{\delta P}(\mathbf{r}_i, E_i, \Omega_i)}{\sum_{\mathcal{T} \in T^{\text{AMS}}} w_{\text{AMS}}(\mathcal{T}) \sum_{P_i \in \mathcal{T}} \mathbf{1}_{\delta P}(\mathbf{r}_i, E_i, \Omega_i)}$$

with AMS weights given by the  
genealogy formula

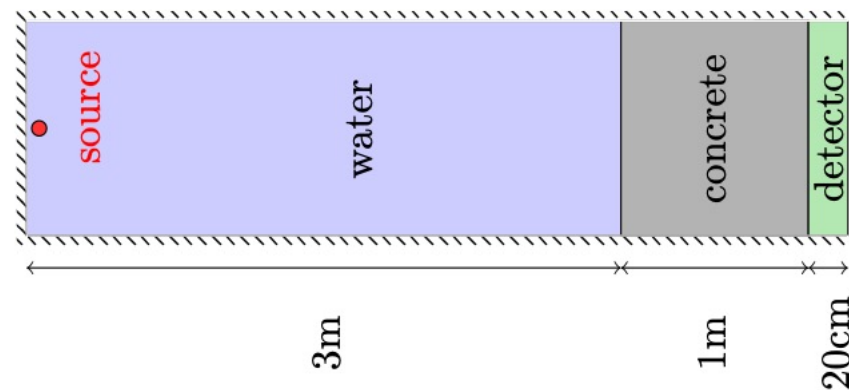
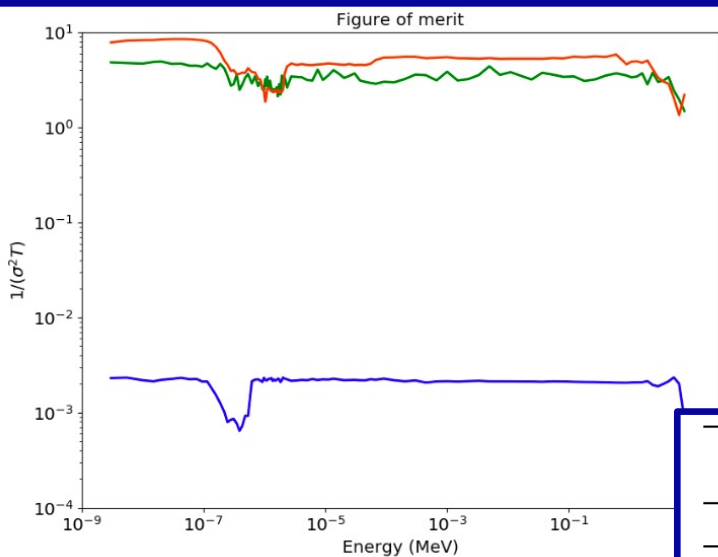


$$w_{\text{AMS}}(\mathcal{T}) = \frac{1}{N} \prod_{i=0}^j \left(1 - \frac{K_i}{N}\right)$$

$$\text{for } \mathcal{T} \in T^{\text{AMS}} = \left( \bigcup_{q=0}^Q S_q^{\text{off}} \right) \cup S_Q^{\text{on}}$$

AMS level  $\longrightarrow$

importance  $\longrightarrow$



| Importance map<br>(properties)             | Mean( $n/cm^3/s$ )     | error(%) | time(s)            | FOM<br>(arbitrary units) |
|--|------------------------|----------|--------------------|--------------------------|
| <b>Adaptive Multilevel Splitting</b>       |                        |          |                    |                          |
| INIPOND<br>(Manual mode)                   | $2.58 \times 10^{-15}$ | 9.90     | $1.67 \times 10^5$ | 9                        |
| IDT<br>(cross sections from TRIPOLI4)      | $2.61 \times 10^{-15}$ | 9.83     | $1.20 \times 10^5$ | 13                       |
| IDT<br>(cross sections from APOLLO3)       | $2.78 \times 10^{-15}$ | 7.11     | $1.59 \times 10^5$ | 16                       |
| <b>Adjoint score<br/>(scored with AMS)</b> | $2.66 \times 10^{-15}$ | 9.98     | $1.08 \times 10^5$ | 13                       |
| <b>Exponential Transform</b>               |                        |          |                    |                          |
| INIPOND<br>(Manual mode)                   | $2.55 \times 10^{-15}$ | 6.51     | $9.41 \times 10^4$ | 39                       |
| IDT<br>(cross sections from TRIPOLI4)      | $2.04 \times 10^{-15}$ | 6.60     | $2.39 \times 10^5$ | 23                       |
| IDT<br>(cross sections from APOLLO3)       | $2.81 \times 10^{-15}$ | 0.82     | $3.27 \times 10^3$ | 576                      |
| <b>Adjoint score<br/>(scored with AMS)</b> | $2.77 \times 10^{-15}$ | 0.52     | $4.33 \times 10^3$ | 1313                     |

- **AMS with AMS  $\xi$  has a FOM equivalent to AMS with deterministic  $\xi$**
- **ET with AMS  $\xi$  has FOM 2 times than ET with deterministic  $\xi$**

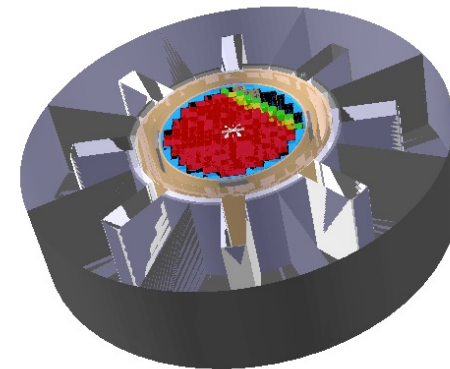


# □ Use of AMS in reactor physics

- Chain reaction & population control
- Spatial correlations
- AMS & branchless collisions
- Results

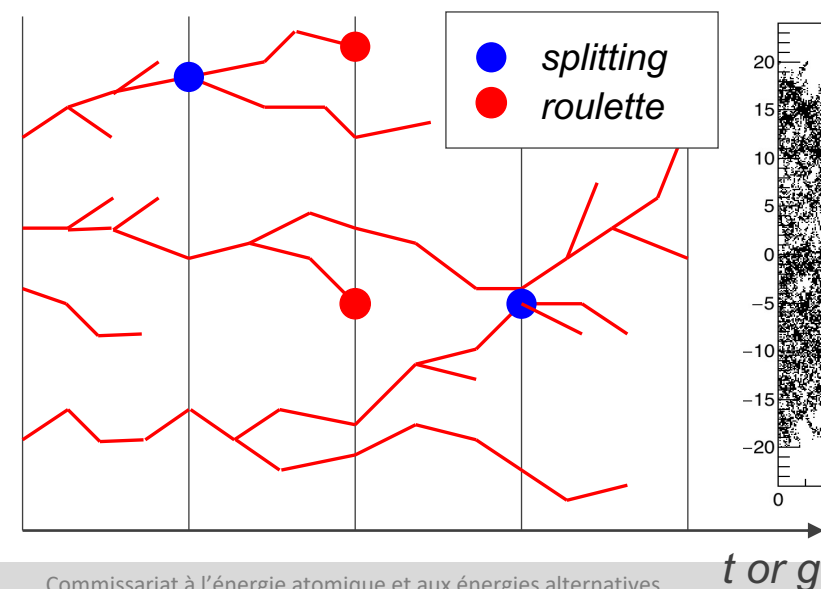


- ❑ Neutron transport in fissile media (**birth/death-killing process**)
- ❑ Critical Boltzmann equation
- ❑ Simplified mono-E model : **BBBM with population control**
- ❑ Population control usually done via '**power iteration**'
  
- ❑ Eigenvalue: **reproduction factor  $\alpha$  (t) or keff (g)**
- ❑ Eigenvectors: **power distribution**

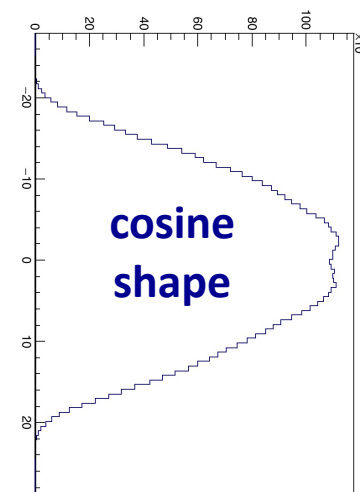
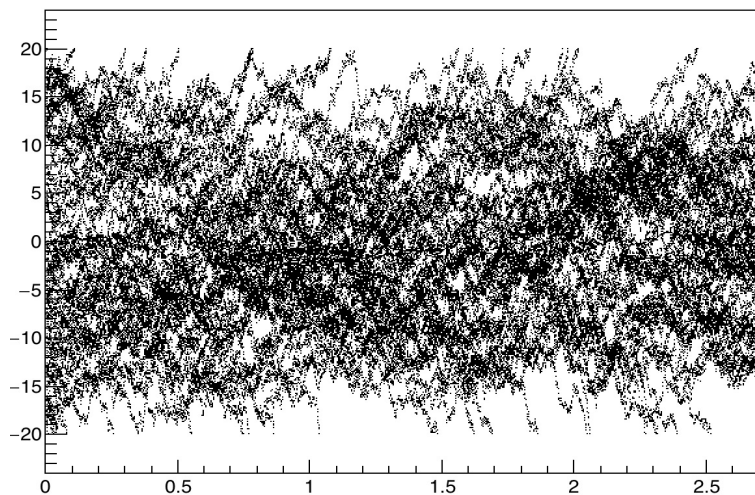


900 MW PWR

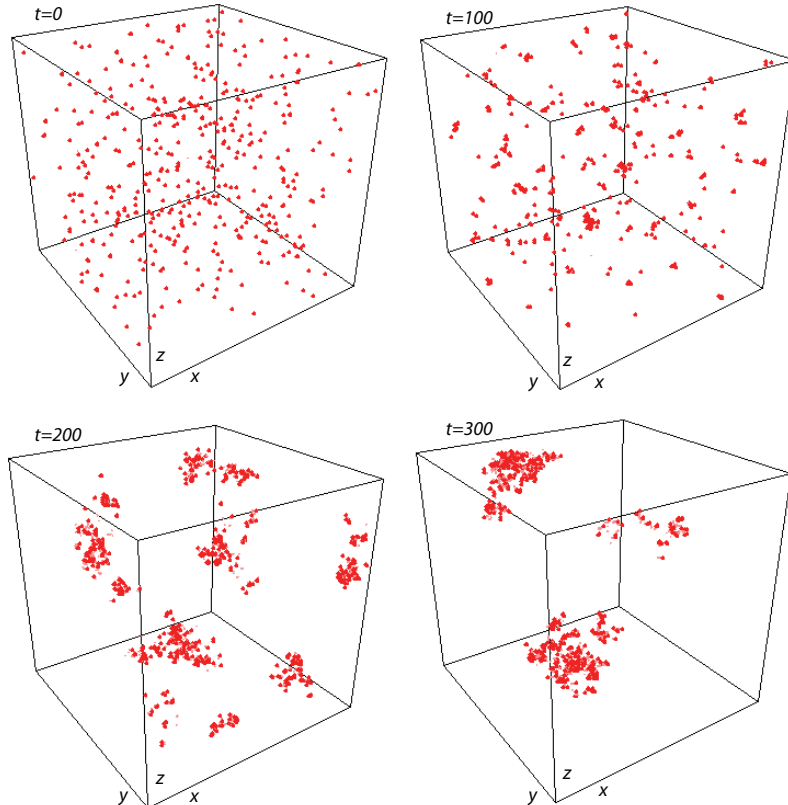
Population control algo. to keep  $N$  constant



1D mono-E rod

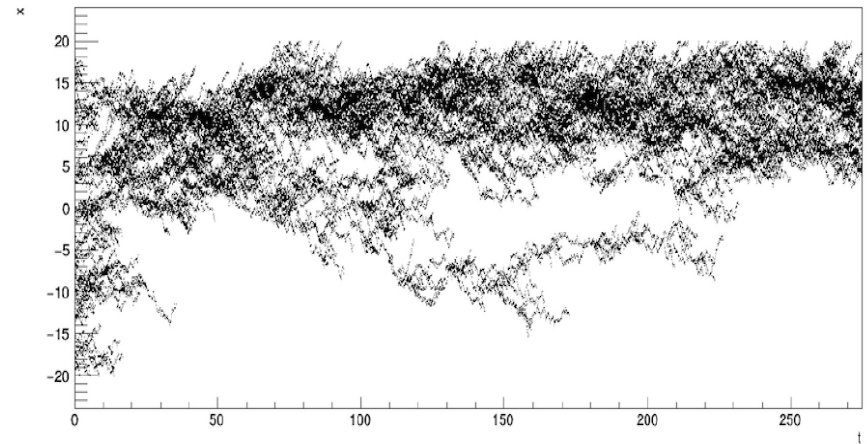


- Branching nature of the **fission process** induces spatial correlations & clustering

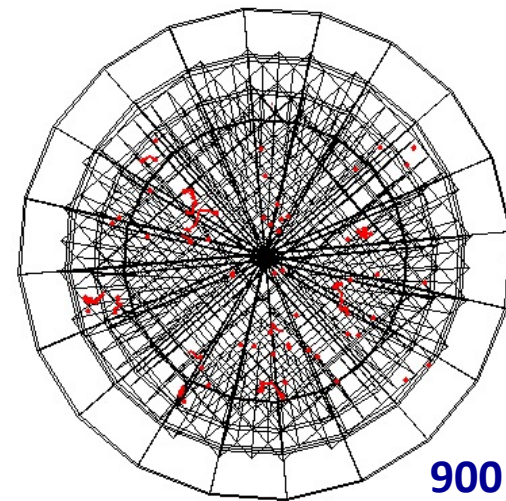


$$g_t(r) = \frac{\lambda v_2}{8Dc_0\pi^{3/2}r} \Gamma\left(\frac{1}{2}, \frac{r^2}{8Dt}\right)$$

(Dumonteil, 2014)



**decoupled 1D mono-E rod**



**900 MW PWR**

- Power iteration** induces a saturation of spatial correlations **while preserving them**

- **AMS** can be seen as a tool for both
  - **population control**
  - **variance reduction**
- **Example:**
  - $k_{eff} < 1$
  - **detector = time/generation**
  - **rare event = surviving population**
- **Similarity with a Fleming-Viot particle system**

number of neutrons to kill

number of iterations to reach T

target time

$$\left(1 - \frac{\pi k_i}{N}\right)^I = e^{-d \cdot T}$$

number of initial AMS neutrons

if  $d < 0 \Rightarrow$  neutron survival probability at T

$$\hat{d} = -\frac{I}{T} \ln\left(1 - \frac{\pi k_i}{N}\right)$$

number of neutrons to kill

number of iterations to reach  $g_0$

target generation

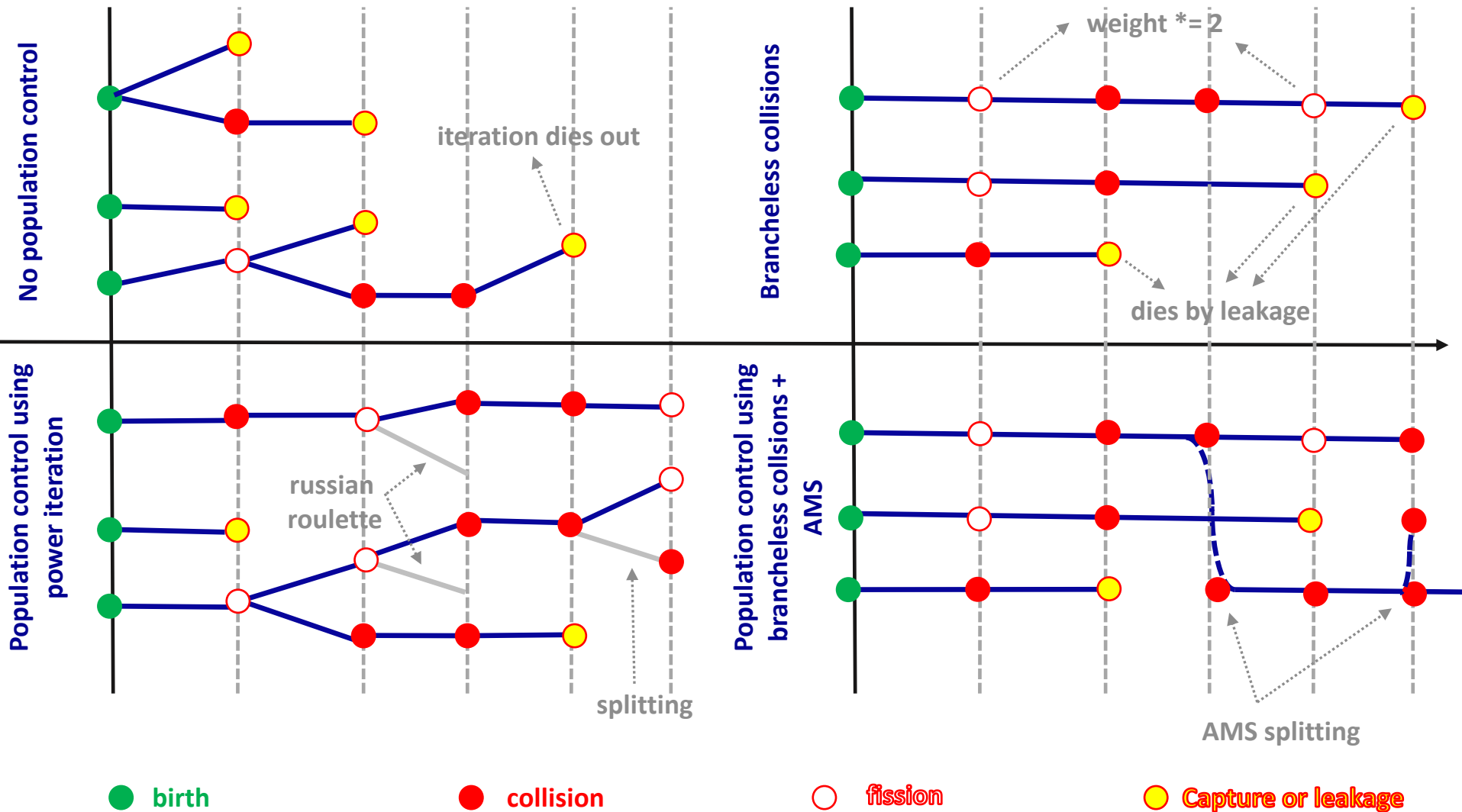
$$\left(1 - \frac{\pi k_i}{N}\right)^I = k_{eff}^{g_0}$$

number of initial AMS neutrons

if  $k_{eff} < 1 \Rightarrow$  neutron survival probability at  $g_0$

$$\hat{k}_{eff} = e^{-\frac{I \pi k_i}{g_0 N}}$$

AMS used in combination with **branchless collisions**

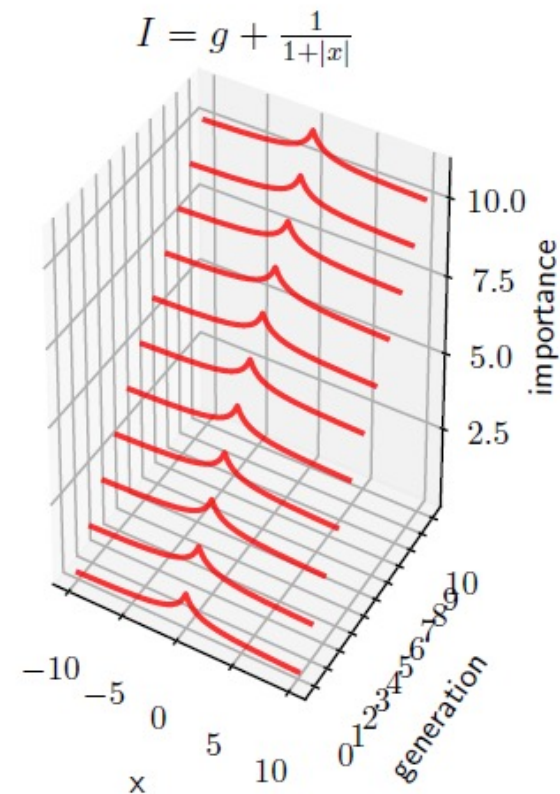


- ❑ « Robustness » of AMS : tracks only need to be ranked (absolute value of importance has no physical meaning)

$$\Rightarrow \text{Importance} = g + \frac{1}{1 + |x|}$$

⇒ **Dominant term to rank tracks by generations & push neutrons through generations**

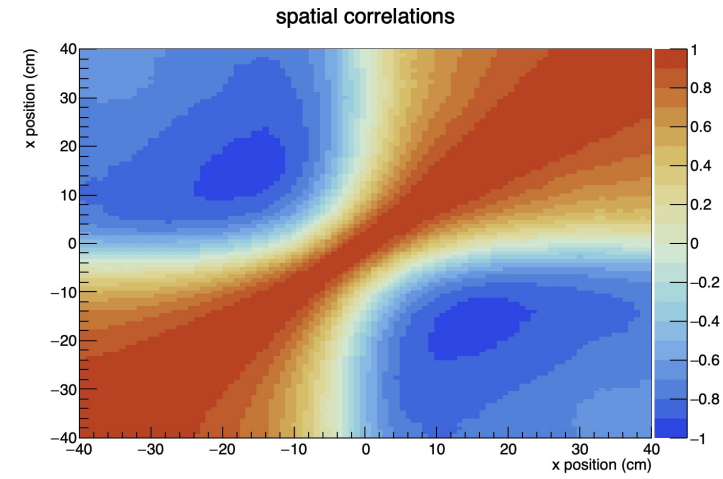
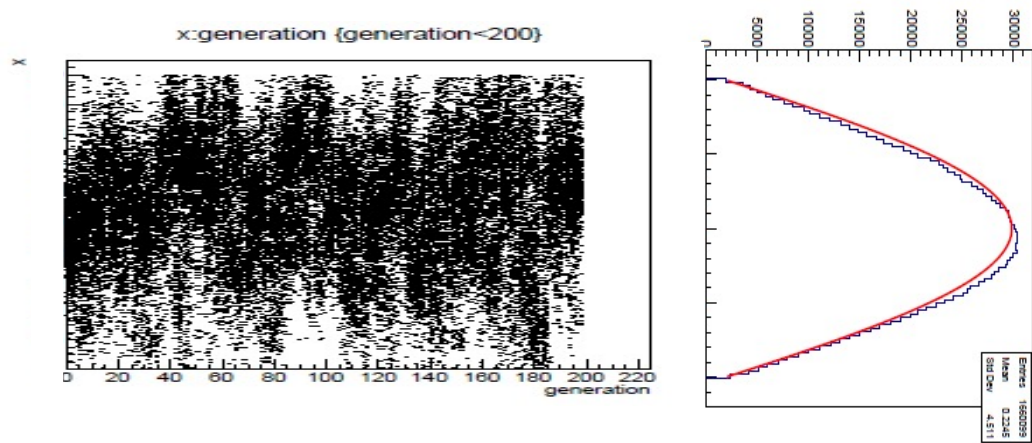
|  $< 1$  : Ensures to rank neutrons inside a generation  
 | A purely discrete importance can cause the algorithm to prematurely stop



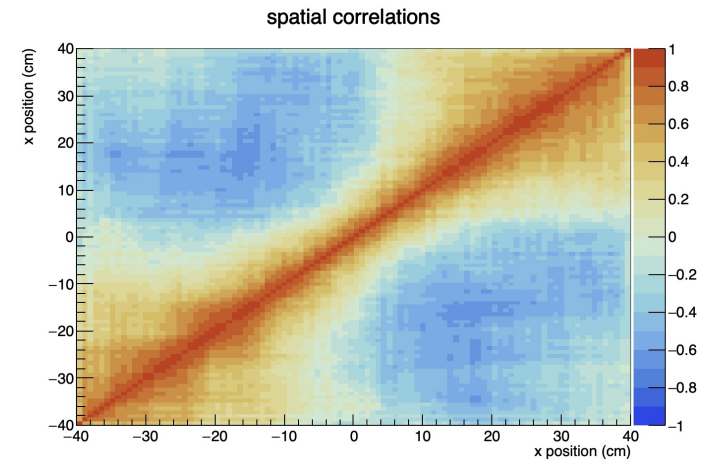
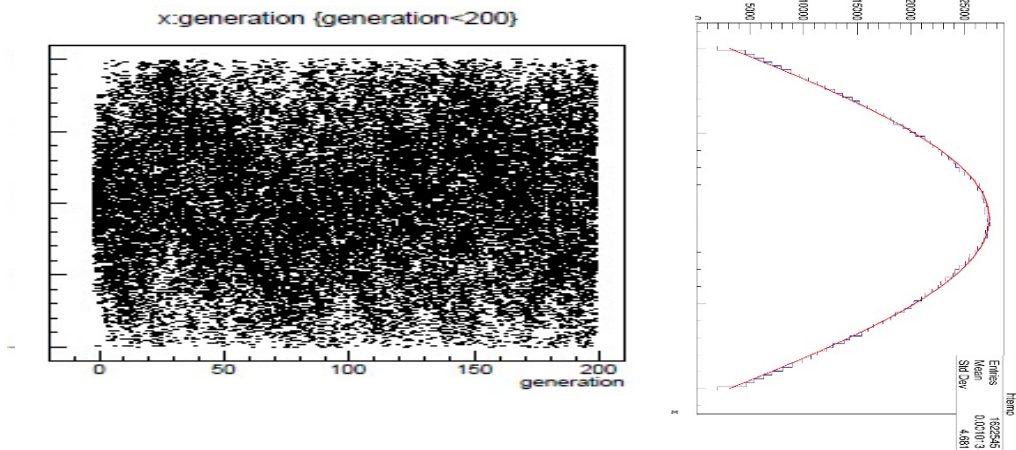
(Phd Kevin Frohlicher, 2022)

- ❑ 80 cm slab / binary branching 'almost'-Brownian motion
- ❑ 100 independent simulations / 1000 neutrons per generation / 1000 generations
- ❑ Spatial correlations are **strongly attenuated (less clustering)**

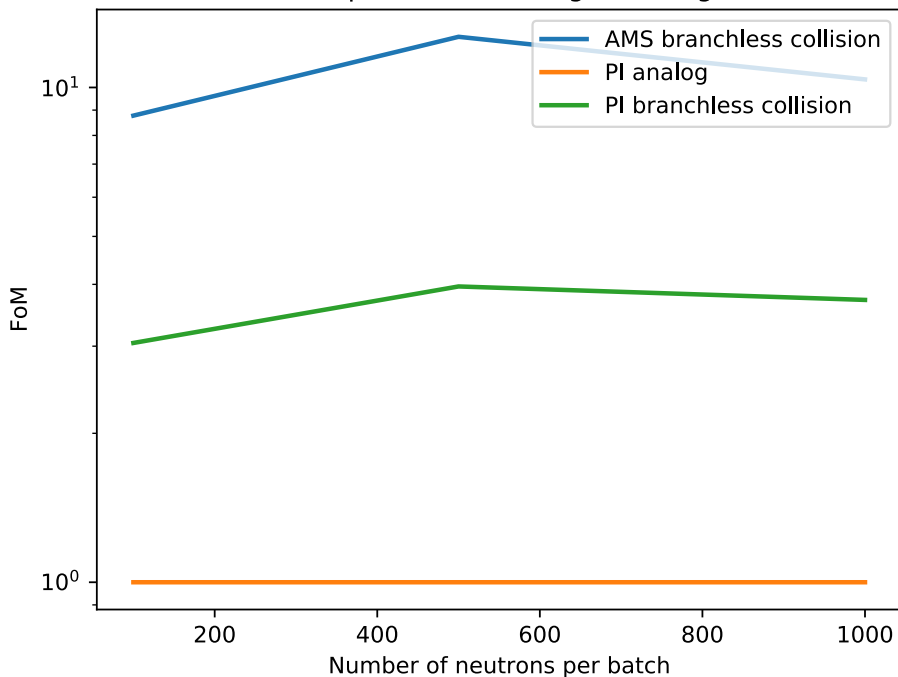
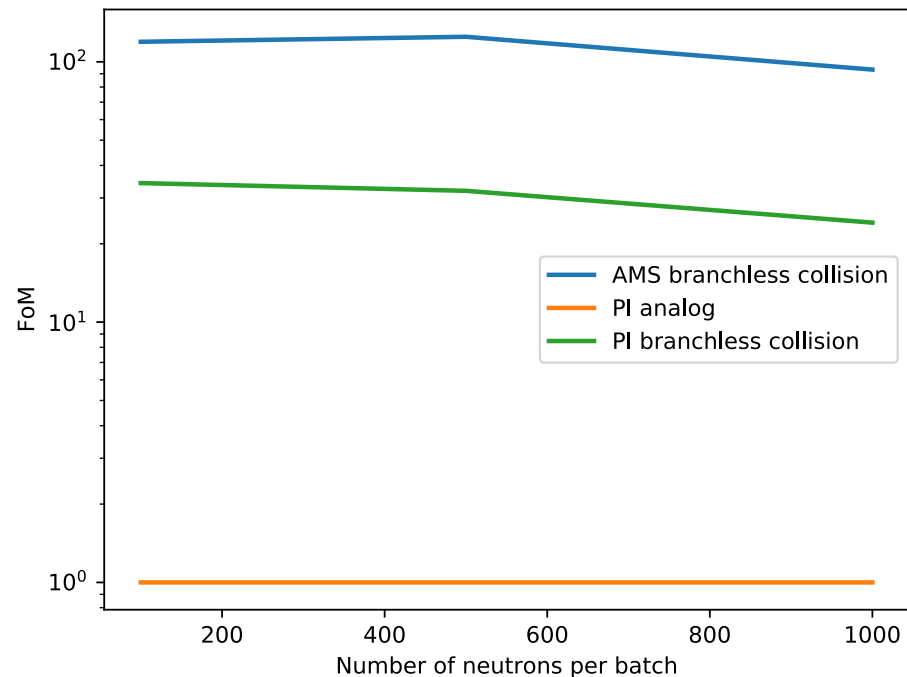
Power iteration  
analog



AMS branchless  
collision



Mean FoM for spatial flux averaged over generations

FoM for the  $k_{eff}$  averaged over generations

- ❑ FOMs are sensitively improved
- ❑ And spatial correlations are tempered
- ❑ Leaves room for (spatial,directional,energetic,generational/time) variance reduction !

- ❑ **AMS vs exponential transform :**
  - If exponential transform uses 0-variance schemes, AMS exhibits **lower FOM**
  - Way **more robust** (the cost function is only sensitive to ranking)
  - Requires **less specific user skills** (0-variance schemes of ET can take weeks to be tuned)
  - Only viable option to **preserve correlations**
  
- ❑ **Already used in “production code” by the nuclear industry** (TRIPOLI-4<sup>®</sup>) in radiation protection contexts
  
- ❑ **Developments on their way to popularize the method in (astro-)particle physics** (Geant4)
  - Physics beyond standard model. Ex: background calculation of elastic neutrino-nucleus scattering experiments
  - Dark matter experiments where both signal&background look for rare events
  
- ❑ **Openings towards quantum mechanics codes (diffusion Monte Carlo)** following the developments of AMS for eigenvalue/eigenvectors problems



- ❑ F. Cérou and A. Guyader, “Adaptive Multilevel Splitting for Rare Event Analysis”, *Stochastic Anal.Appl.* 25,2, 417 (2007).
- ❑ F. Cérou, et al. “A multiple replica approach to simulate reactive trajectories”, *Journal of Chemical Physics*, 134, 054108, (2011).
- ❑ D. Aristoff, T. Lelièvre, C.G. Mayne and I. Teo. “Adaptive multilevel splitting in molecular dynamics simulations” *ESAIM Proc.*, 48:215-225, (2015).
- ❑ CE Bréhier, M. Gazeau, L. Goudenège, et al. “Unbiasedness of some generalized Adaptive Multilevel Splitting algorithms.” *The Annals of Applied Probability*, 26(6):3559– 3601 (2016).
  
- ❑ H. Louvin et al., “Adaptive Multilevel Splitting for MonteCarlo Particle Transport”, *Eur. Phys. J. Nucl. Sci. Technol.* 3,29 (2017).
- ❑ H. Louvin, “Development of an Adaptive VarianceReduction Technique for Monte Carlo Particle Transport”, PhD Thesis, Paris-Saclay University (2017).
- ❑ M. Nowak et al, “Accelerating Monte Carlo Shielding Calculations inTRIPOLI-4 with a Deterministic Adjoint Flux », *Nucl. Sci. Eng.*, 193 (2019).
- ❑ Mancusi et al., « Evaluating importance maps for using deterministic or on-line methods », *ANS RPSD Proceedings* (2018)

- ❑ M. Nowak et al, “Accelerating Monte Carlo Shielding Calculations in TRIPOLI-4 with a Deterministic Adjoint Flux », PhD thesis, Paris-Saclay University (2018).
- ❑ E. Brun et al, “Tripoli-4<sup>®</sup>, CEA, EDF and AREVA Reference Monte Carlo Code”, Ann. Nucl. Energy 82,151 (2015).
- ❑ S. Agostinelli et al., “Geant4, a simulation toolkit”, Nuclear Instruments and Methods in Physics A, Vol 506 (2003)
- ❑ I. ZMIJAREVIC et al, “Deterministic Solutions for 3D Kobayashi Benchmarks,” Prog. Nucl. Energy, 39, 207 (2001); [https://doi.org/10.1016/S0149-1970\(01\)00013-0](https://doi.org/10.1016/S0149-1970(01)00013-0).
- ❑ D. SCHNEIDER et al., “Apollo<sup>®</sup>: CEA/DEN Deterministic Multi-Purpose Code for Reactor Physics Analysis,” Proc. PHYSOR 2016 Conf., Sun Valley, Idaho, May 1–5, 2016, p. 2274, American Nuclear Society (2016).



**Thank you for your attention**